



# Room-temperature ferromagnetism and ferroelectricity of nanocrystalline $\text{La}_2\text{Ti}_2\text{O}_7$

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## ABSTRACT

$\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders and pellets were prepared by sol–gel method. The nanocrystalline powders present room-temperature ferromagnetism (FM). The vacuum annealing enhances FM, showing that the observed room-temperature FM for  $\text{La}_2\text{Ti}_2\text{O}_7$  possibly originates from oxygen vacancies at/near surfaces of nanograins.  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets annealed at 1000 °C for 1 h show the coexistence of ferromagnetism and ferroelectricity (FE) at room temperature. Meanwhile, the room-temperature magnetodielectric (MD) effect was also observed in the  $\text{La}_2\text{Ti}_2\text{O}_7$  multiferroic pellet, indicating the coupling between magnetic and electric properties. The value of  $\Delta\epsilon_r/\epsilon_r(0)$  for  $\text{La}_2\text{Ti}_2\text{O}_7$  multiferroic pellet reaches 39.3% at 1 kHz under  $H = 6$  kOe.

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## 1. Introduction

The multiferroic materials with coexistence of ferromagnetism (FM) and ferroelectricity (FE) are of particular interest since they present possibility for potential applications in spintronics, information storage and sensors [1]. Generally, most ferromagnetic oxides contain the center of symmetry and do not show electric polarization, while most sintered ferroelectric bulks are not magnetic [2]. However, more and more multiferroic materials were reported in recent years [3–13], e.g.  $\text{TbMnO}_3$  [3,7],  $\text{BiFeO}_3$  [5] and  $\text{BiMnO}_3$  [6]. It can be expected that for the multiferroics with ferromagnetic and ferroelectric ordering simultaneously, the coupling between the magnetic and electric properties will lead to the magnetoelectric (ME) effect [1,3,5–7,9] in which the magnetization can be controlled by application of electric fields, and vice versa. The coupling can also give rise to the magnetodielectric (MD) [11,14] or magnetocapacitance (MC) effects [6,10], in which dielectric permittivity or capacitance can be adjusted by an applied magnetic field. Multiferroic composite materials consisting of both ferroelectric and ferromagnetic phases could yield giant magnetoelectric coupling response above room temperature, which are of potential applications [15,16].

Undoped oxides, such as  $\text{HfO}_2$ ,  $\text{TiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{In}_2\text{O}_3$ ,  $\text{ZnO}$  and  $\text{MgO}$ , have been found to exhibit weak FM at room temperature [17–24]. The observed FM may be related to the structural defects and size-confinement effects. Point defects in insulators might create

localized electronic levels within the band gap, leading to different charge states and magnetic moments [25]. There were experimental evidences that the FM in  $\text{HfO}_2$  and  $\text{TiO}_2$  powders/films originated from oxygen vacancies [22,23]. However, an *ab initio* calculation demonstrated that cation vacancies were responsible for the magnetic moments in  $\text{HfO}_2$  undoped oxide [26]. Both experimental and *ab initio* calculation [24,27] results showed that FM in nanocrystalline  $\text{MgO}$  originated from Mg vacancies at/near surfaces of nanograins.

Perovskite  $\text{BaTiO}_3$  is a traditional ferroelectric material, and  $\text{BaTiO}_3$  bulks sintered at high temperature are nonmagnetic. Recently, Mangalam et al. [28] reported the multiferroic properties with the coexistence of FM and FE in  $\text{BaTiO}_3$  nanocrystalline (300 nm) pellet annealed at 1000 °C for 1 h. By first-principle calculation, they suggested that the FM of nanocrystalline  $\text{BaTiO}_3$  might come from the oxygen vacancies and ferroelectricity from the core. A magnetocapacitance effect was observed in  $\text{BaTiO}_3$  nanocrystalline (300 nm) pellet at room temperature [28].  $\text{La}_2\text{Ti}_2\text{O}_7$ , which is a ferroelectric material with a high ferroelectric Curie temperature of 1500 °C [29,30], is a promising candidate for high-temperature piezoelectric and electro-optic devices [31–33]. Since the magnetic moments of both La atom and  $\text{La}^{3+}$  are zero and Ti element is nonmagnetic [34], there was no interest in the magnetism in bulk  $\text{La}_2\text{Ti}_2\text{O}_7$  before. Whether the nanocrystalline  $\text{La}_2\text{Ti}_2\text{O}_7$  is ferromagnetic or not? In this letter, we show that similar to undoped oxides  $\text{HfO}_2$ ,  $\text{TiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{In}_2\text{O}_3$ ,  $\text{ZnO}$ ,  $\text{MgO}$  and  $\text{BaTiO}_3$  [17–24,28],  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders present room-temperature FM. The vacuum annealing enhances FM, indicating that the FM in nanocrystalline  $\text{La}_2\text{Ti}_2\text{O}_7$  possibly originates from oxygen vacancies at/near the surfaces of nanograins. The coexistence of FM

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**Table 1**  
Different types of  $\text{La}_2\text{Ti}_2\text{O}_7$  samples.

$T_A$ ( $^\circ\text{C}$ )	Annealing conditions		
700 (powders)	2 h in air	2 h in air + 20 min in vacuum	2 h in air + 20 min in vacuum + 20 min in air
800 (powders)	2 h in air		
1000 (powders)	2 h in air		
1000 (pellets)	30 min in air	1 h in air	1.5 h in air
1100 (pellet)	1 h in air		

and FE at room-temperature could be observed in  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets annealed at appropriate low temperatures and in proper time period. Meanwhile, the MD effect was also observed, which implied the coupling between FM and FE.

## 2. Experimental

$\text{La}_2\text{Ti}_2\text{O}_7$  powders were prepared by sol-gel method. 0.2 mol  $\text{La}(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$  with a purity of 99.99 wt% (without any ferromagnetic elements) and 0.2 mol tetra-*n*-butyl titanate with a purity of 99.99 wt% (without any ferromagnetic elements) were dissolved in 180 mL ethanol, with constant stirring at about  $35^\circ\text{C}$ , to get the sol. The gel was obtained by stirring the sol at  $65^\circ\text{C}$  for two days. The precursor was annealed at  $700^\circ\text{C}$ ,  $800^\circ\text{C}$  and  $1000^\circ\text{C}$  for 2 h in air, respectively. Parts of powders after annealing at  $700^\circ\text{C}$  were reheated in vacuum at  $700^\circ\text{C}$  for 20 min, and then parts of vacuum annealed powders were backfired in air at  $700^\circ\text{C}$  for 20 min. Some  $\text{La}_2\text{Ti}_2\text{O}_7$  powders were pressed with a mold (made from nylon), and then annealed at  $1000^\circ\text{C}$  and  $1100^\circ\text{C}$ , respectively. All the different types of  $\text{La}_2\text{Ti}_2\text{O}_7$  samples are listed in Table 1. The structures of  $\text{La}_2\text{Ti}_2\text{O}_7$  powders annealed at different temperatures were investigated by X-ray diffraction with  $\text{Cu K}\alpha$  radiation. The microstructure of the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h was observed by field emission scanning electron microscope (FE-SEM). The magnetic properties of all the samples were measured using alternating gradient magnetometer (AGM) at room-temperature. The ferroelectric properties of  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets were measured by a ferroelectric meter TF Analyzer 2000. The MD effect was measured with a HP 4294A impedance analyzer. Both surfaces of the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets for ferroelectric and MD measurements were deposited with Ag.

## 3. Results and discussion

The X-ray diffraction patterns of  $\text{La}_2\text{Ti}_2\text{O}_7$  as-synthesized powders are shown in Fig. 1. At  $700^\circ\text{C}$ ,  $\text{La}_2\text{Ti}_2\text{O}_7$  powders are partly crystallized with an orthorhombic structure (group  $Pna21$ ). The powders annealed above  $800^\circ\text{C}$  crystallize with a monoclinic structure (group  $P2_1$ ). Such phase transformation has been reported previously [35]. The average grain sizes of  $\text{La}_2\text{Ti}_2\text{O}_7$  powders were estimated to be 28 nm for  $700^\circ\text{C}$  and 33 nm for  $800^\circ\text{C}$  by Scherrer's method.

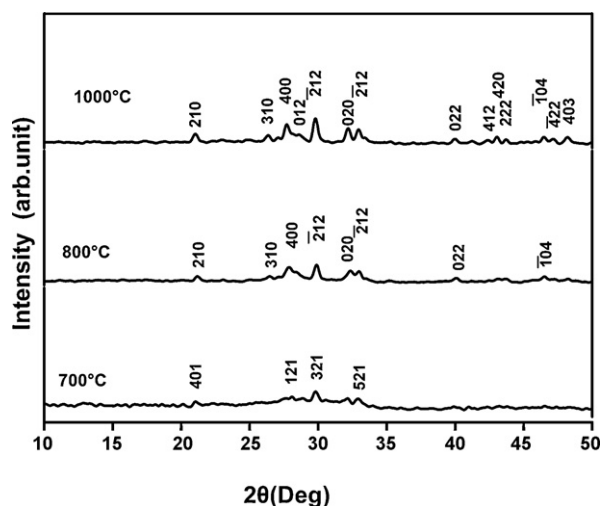


Fig. 1. The X-ray diffraction patterns of  $\text{La}_2\text{Ti}_2\text{O}_7$  powders.

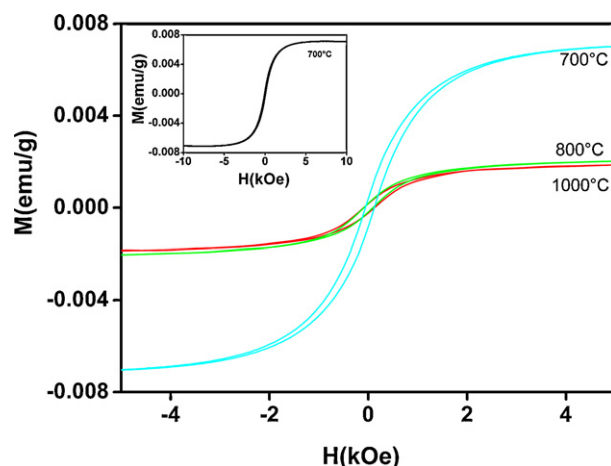


Fig. 2. The room-temperature  $M$ - $H$  curves after DM corrections for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders with different annealing temperature from  $700$  to  $1000^\circ\text{C}$  for 2 h. The inset is the room-temperature  $M$ - $H$  curve (after DM correction) in a large field range for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders with  $T_A = 700^\circ\text{C}$ .

The magnetization ( $M$ ) versus magnetic field ( $H$ ) curves after diamagnetism (DM) correction for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders at room-temperature were shown in Fig. 2.  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders present hysteresis loops, showing room-temperature FM. The inset is the room-temperature  $M$ - $H$  curve (after DM correction) in a large field range for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders with  $T_A = 700^\circ\text{C}$ . The values of  $M_s$  are  $7.1 \times 10^{-3}$  emu/g for  $T_A = 700^\circ\text{C}$ ,  $2.1 \times 10^{-3}$  emu/g for  $T_A = 800^\circ\text{C}$ , and  $1.9 \times 10^{-3}$  emu/g for  $T_A = 1000^\circ\text{C}$ . In general, the room-temperature saturation magnetization  $M_s$  of oxide nanopowders, which contain no ferromagnetic element, is very small. For example, the value of  $M_s$  is  $\sim 10^{-3}$  emu/g for  $\text{BaTiO}_3$  [28]. The FM of  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders reduces with the increase of annealing temperatures  $T_A$  from  $700$  to  $1000^\circ\text{C}$ , possibly due to the decrease of concentration of structural defects at/near the surfaces of grains.

Fig. 3 shows room-temperature  $M$ - $H$  curves for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders annealed at  $1000^\circ\text{C}$  (after DM correction), and pellets annealed at  $1000^\circ\text{C}$  for 30 min (after DM correction), 1 h (after DM correction), 1.5 h (without DM correction) and  $1100^\circ\text{C}$  for 1 h (without DM correction). The  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets annealed at  $T_A = 1000^\circ\text{C}$  for 30 min and 1 h also present room-temperature FM. A typi-

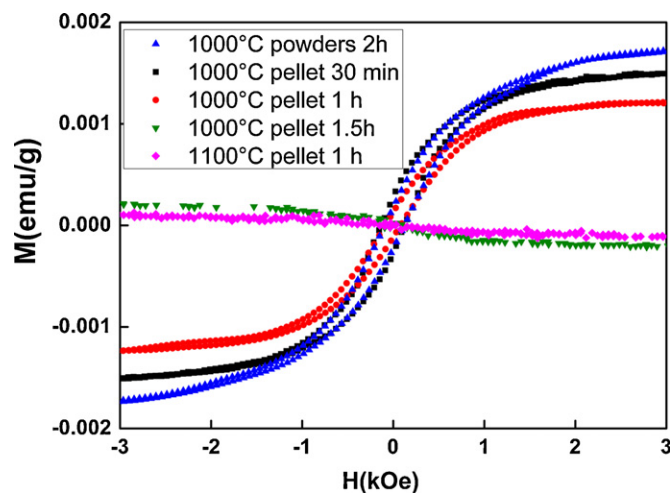


Fig. 3. The room-temperature  $M$ - $H$  curves for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders annealed at  $1000^\circ\text{C}$  (after DM correction), and pellets annealed at  $1000^\circ\text{C}$  for 30 min (after DM correction), 1 h (after DM correction), 1.5 h (without DM correction) and  $1100^\circ\text{C}$  for 1 h (without DM correction).

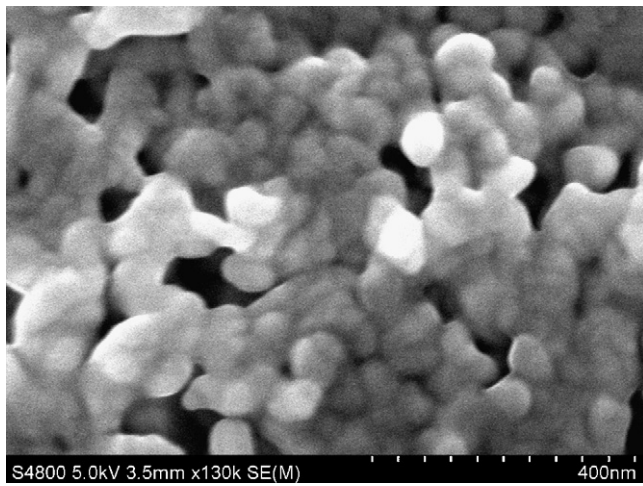


Fig. 4. The FE-SEM image of  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h.

cal microstructure observed by a FE-SEM for the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h was shown in Fig. 4. The volume fraction of grains in this pellet is 15% for 20–50 nm grains, 19% for 50–100 nm grains, 66% for 100–225 nm grains. The average grain size of  $\text{La}_2\text{Ti}_2\text{O}_7$  conventional bulks sintered at  $1150^\circ\text{C}$  for 7 h was about 10–20  $\mu\text{m}$  [32]. The magnetism of  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets depends sensitively on both annealing temperature and annealing time.  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets annealed at  $T_A = 1000^\circ\text{C}$  for 1.5 h and  $1100^\circ\text{C}$  for 1 h exhibit diamagnetism (see Fig. 4). This is consistent with the fact that conventional  $\text{La}_2\text{Ti}_2\text{O}_7$  sintered bulks are nonmagnetic. The values of  $M_s$  of  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets ( $T_A = 1000^\circ\text{C}$  for 30 min and 1 h) are lower than that of  $\text{La}_2\text{Ti}_2\text{O}_7$  powders ( $T_A = 1000^\circ\text{C}$  for 2 h), due to the decrease of defect concentrations at the grain surfaces.

The observed FM in  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders does not originate from impurities, since the reactants do not contain any other ferromagnetic elements, and La/ $\text{La}^{3+}$ , as well as Ti, are nonmagnetic. Similar to the cases of undoped oxides  $\text{HfO}_2$ ,  $\text{TiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{In}_2\text{O}_3$ ,  $\text{ZnO}$ ,  $\text{MgO}$  and  $\text{BaTiO}_3$  [17–24,28], the observed FM in nanocrystalline  $\text{La}_2\text{Ti}_2\text{O}_7$  should originate from structural defects. By first-principle calculation, Mangalam et al. suggested that the FM of nanocrystalline  $\text{BaTiO}_3$  might come from the oxygen vacancies [28]. In order to experimentally verify whether FM originates from oxygen vacancies or cation vacancies, we have examined the influence of the vacuum annealing on magnetism of nanocrystalline  $\text{La}_2\text{Ti}_2\text{O}_7$ . Usually, the vacuum annealing effectively brings about some oxygen vacancies in semiconducting oxides. We found that after annealing at  $700^\circ\text{C}$  in air, a subsequent annealing in vacuum at  $700^\circ\text{C}$  for 20 min enhanced room-temperature FM of  $\text{La}_2\text{Ti}_2\text{O}_7$  powders. The  $M_s$  value for the  $\text{La}_2\text{Ti}_2\text{O}_7$  powders with the backfire in air after vacuum annealing is lower than that of the sample with vacuum annealing, but higher than that of the sample annealed at air (see Fig. 5). It shows that FM in  $\text{La}_2\text{Ti}_2\text{O}_7$  powders mainly originates from oxygen vacancies. It was also reported that the FM of nanocrystalline  $\text{HfO}_2$ ,  $\text{TiO}_2$  and  $\text{BaTiO}_3$  might come from the oxygen vacancies [21,23,28].

It is well known that  $\text{La}_2\text{Ti}_2\text{O}_7$  sintered bulks present ferroelectric properties with a high ferroelectric Curie temperature and a high coercive field [29,30]. Usually, bulk materials with strong ferroelectric properties are prepared by sintering at high temperature. However, for  $\text{BaTiO}_3$  [28] and  $\text{La}_2\text{Ti}_2\text{O}_7$ , FM properties disappear when the annealing temperature  $T_A$  is too high. In order to obtain the coexistence of FM and FE, the annealing temperature  $T_A$  should be carefully chosen. Fig. 6 shows the room-temperature polarization ( $P$ ) versus electric field ( $E$ ) curves for the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $T_A = 1000^\circ\text{C}$  for 1 h. Hysteresis  $P$ – $E$  loops were clearly

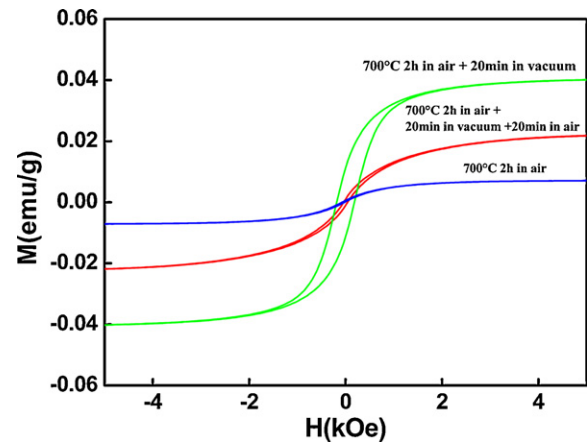


Fig. 5. The room temperature  $M$ – $H$  curves after DM corrections for  $\text{La}_2\text{Ti}_2\text{O}_7$  powders with different annealing conditions at  $700^\circ\text{C}$ .

observed for the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet, showing the room-temperature FE. The  $P$ – $E$  loops are not saturated, since the saturation electrical field is larger than what the equipment used here can provide.

The ferroelectric signal of the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealing at  $T_A = 1000^\circ\text{C}$  for 1 h is weak and the loops are slim (see Fig. 6). These two phenomena are possibly connected with the low annealing temperature and the small applied electric field. The low annealing temperature with short annealing time leads to small-size grains. Many investigators have investigated grain-size effect on ferroelectric properties [36–39]. There is a reduction of FE with the decrease of grain size for ferroelectric nanomaterials. In addition, the  $P$ – $E$  loops are not continuous (see Fig. 6) at the third and fourth quadrants. Similar  $P$ – $E$  loops were also observed by other groups  $\text{Bi}_{1-x}\text{Ba}_x\text{FeO}_3$  ( $x=0.15$ ) [9].

Our experimental results demonstrate that the  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline pellets annealed at  $T_A = 1000^\circ\text{C}$  for 1 h is a room-temperature multiferroic material with coexistence of FM and FE, even though both FM and FE are weak. The coexistence of FM and FE for  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets requires appropriate annealing temperature and time. Fig. 7 shows the AC electric field frequency dependence of dielectric constant ( $\epsilon_r$ ) measured at room-temperature (without DC magnetic field) for the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $T_A = 1000^\circ\text{C}$  for 1 h. It is clear that the dielectric constant decreases with increasing frequency. This result can be explained by the phenomenon of dipole relaxation wherein at low frequencies the dipoles are able to follow the frequency of the applied field [9]. In order

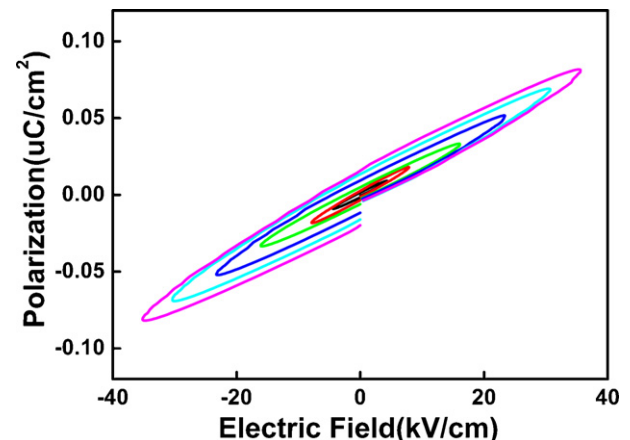


Fig. 6.  $P$ – $E$  curves of  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h.

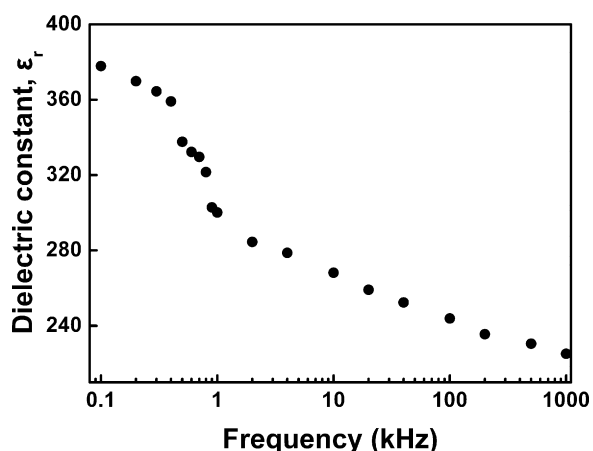


Fig. 7. Applied AC electric field frequency dependence of dielectric constant measured at room-temperature (without DC magnetic field) for  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h.

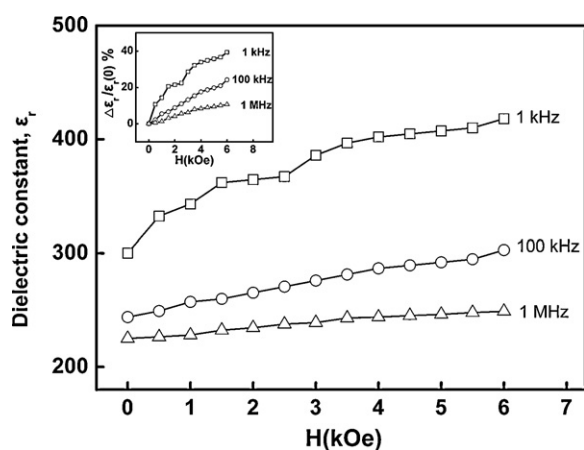


Fig. 8. Applied magnetic field dependence of dielectric constant measured at 1 kHz, 100 kHz and 1 MHz (at room temperature) for  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet annealed at  $1000^\circ\text{C}$  for 1 h. The inset shows the magnetodielectric (MD) effect.

to test the coupling between electric and magnetic properties in the room-temperature multiferroic  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet, the applied magnetic field dependence of dielectric constant was measured at various frequencies (see Fig. 8). Room-temperature magnetodielectric effect was clearly observed. The field dependence of  $\Delta\varepsilon_r/\varepsilon_r(0) = [\varepsilon_r(H) - \varepsilon_r(0)]/\varepsilon_r(0)$  for the  $\text{La}_2\text{Ti}_2\text{O}_7$  pellet was also plotted in the inset of Fig. 8. With increasing magnetic field, the dielectric constant increases. Such results indicated that there is a coupling between magnetic and electric properties. As shown in the inset of Fig. 8, a positive MD effect occurs. The value of  $\Delta\varepsilon_r/\varepsilon_r(0)$  reaches 39.3% at 1 kHz under  $H = 6$  kOe. Positive magnetocapacitance (or magnetodielectric) effects were also observed in  $\text{BaTiO}_3$  nanocrystalline pellets (annealed at  $1000^\circ\text{C}$  for 1 h) [28] and  $\text{Bi}_{1-x}\text{Ba}_x\text{FeO}_3$  compounds [9].

#### 4. Conclusion

$\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders prepared by sol-gel method present room-temperature FM. The FM of nanocrystalline powders may occur at the surface of nanograins. Vacuum annealing enhances FM of  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders. The observed FM in  $\text{La}_2\text{Ti}_2\text{O}_7$  nanocrystalline powders may originate from oxygen vacancies at/near the surfaces of nanograins.  $\text{La}_2\text{Ti}_2\text{O}_7$  pellets annealed at  $T_A = 1000^\circ\text{C}$  for 1 h present both room-temperature FM and FE. Meanwhile,  $\text{La}_2\text{Ti}_2\text{O}_7$  multiferroic pellets also exhibit

the room-temperature MD effect, indicating the coupling between magnetic and electric properties.

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#### References

- [1] N.A. Spaldin, M. Fiebig, *Science* 309 (2005) 391.
- [2] N.A. Hill, *J. Phys. Chem. B* 104 (2000) 6694.
- [3] T. Kimura, T. Goto, H. Shintani, K. Ishizaka, T. Arima, Y. Tokura, *Nature* 426 (2003) 55.
- [4] N. Hur, S. Park, P.A. Sharma, J.S. Ahn, S. Guha, S.-W. Cheong, *Nature* 429 (2004) 392.
- [5] J. Wang, J.B. Neaton, H. Zheng, V. Nagarajan, S.B. Ogale, B. Liu, D. Viehland, V. Vaithyanathan, D.G. Schlom, U.V. Waghmare, N.A. Spaldin, K.M. Rabe, M. Wuttig, R. Ramesh, *Science* 299 (2003) 1719.
- [6] T. Kimura, S. Kawamoto, I. Yamada, M. Azuma, M. Takano, Y. Tokura, *Phys. Rev. B* 67 (2003) 180401(R).
- [7] N. Abe, K. Taniguchi, S. Ohtani, H. Umetsu, T. Arima, *Phys. Rev. B* 80 (2009) 020402(R).
- [8] J. Wen, G. Xu, G. Gu, S.M. Shapiro, *Phys. Rev. B* 80 (2009) 020403(R).
- [9] D.H. Wang, W.C. Goh, M. Ning, C.K. Ong, *Appl. Phys. Lett.* 88 (2006) 212907.
- [10] M.P. Singh, W. Prellier, Ch. Simon, B. Raveau, *Appl. Phys. Lett.* 87 (2005) 022505.
- [11] R. Ranjith, Asish K. Kundu, M. Filippi, B. Kundys, W. Prellier, B. Raveau, J. Laverdière, M.P. Singh, S. Jandl, *Appl. Phys. Lett.* 92 (2008) 062909.
- [12] Q.H. Jiang, F.T. Liu, C.-W. Nan, Y.H. Lin, M.J. Reece, H.X. Yan, H.P. Ning, Z.J. Shen, *Appl. Phys. Lett.* 95 (2009) 012909.
- [13] R. Seshadri, N.A. Hill, *Chem. Mater.* 13 (2001) 2892.
- [14] N. Hur, S. Park, P.A. Sharma, S. Guha, S.-W. Cheong, *Phys. Rev. Lett.* 93 (2004) 017207.
- [15] C.W. Nan, M.I. Bichurin, S.X. Dong, D. Viehland, G. Srinivasan, *J. Appl. Phys.* 103 (2008) 031101.
- [16] H. Zheng, J. Wang, S.E. Lofland, Z. Ma, L. Mohaddes-Ardabili, T. Zhao, L. Salamanca-Riba, S.R. Shinde, S.B. Ogale, F. Bai, D. Viehland, Y. Jia, D.G. Schlom, M. Wuttig, A. Roytburd, R. Ramesh, *Science* 303 (2004) 661.
- [17] M. Venkatesan, C.B. Fitzgerald, J.M.D. Coey, *Nature* 430 (2004) 630.
- [18] N.H. Hong, J. Sakai, N. Poirot, V. Brizé, *Phys. Rev. B* 73 (2006) 132404.
- [19] S.D. Yoon, Y. Chen, A. Yang, T.L. Goodrich, X. Zuo, K. Ziemer, C. Vittoria, V.G. Harris, *J. Magn. Magn. Mater.* 309 (2007) 171.
- [20] A. Sundaresan, R. Bhargavi, N. Rangarajan, U. Siddesh, C.N.R. Rao, *Phys. Rev. B* 74 (2006) 161306.
- [21] N.H. Hong, J. Sakai, N.T. Huong, N. Poirot, A. Ruyter, *Phys. Rev. B* 72 (2005) 045336.
- [22] J.M.D. Coey, M. Venkatesan, P. Stamenov, C.B. Fitzgerald, L.S. Dorneles, *Phys. Rev. B* 72 (2005) 024450.
- [23] N.H. Hong, J. Sakai, N.T. Huong, A. Ruyter, V. Brizé, *Appl. Phys. Lett.* 89 (2006) 252504.
- [24] J.F. Hu, Z.L. Zhang, M. Zhao, H.W. Qin, M.H. Jiang, *Appl. Phys. Lett.* 93 (2008) 192503.
- [25] A.M. Stoneham, *Theory of Defects in Solids*, Oxford University Press, Oxford, 1975.
- [26] C.D. Penmaraju, S. Sanvito, *Phys. Rev. Lett.* 94 (2005) 217205.
- [27] F. Gao, J.F. Hu, C.L. Yang, Y.J. Zheng, H.W. Qin, L. Sun, X.W. Kong, M.H. Jiang, *Solid State Commun.* 149 (2009) 855.
- [28] R.V.K. Mangalam, Ray Nirat, Umesh V. Waghmare, A. Sundaresan, C.N.R. Rao, *Solid State Commun.* 149 (2009) 1.
- [29] S. Nanamatsu, M. Kimura, K. Doi, S. Matsushita, N. Yamada, *Ferroelectrics* 8 (1974) 511.
- [30] M. Kimura, S. Nanamatsu, T. Kawamura, S. Matsushita, *Jpn. J. Appl. Phys.* 13 (1974) 147.
- [31] J.K. Yamamoto, A.S. Bhalla, *J. Appl. Phys.* 70 (1991) 4469.
- [32] P.A. Fuieler, R.E. Newnham, *J. Am. Ceram. Soc.* 74 (1991) 2876.
- [33] S. Marzullo, E.N. Bunting, *J. Am. Ceram. Soc.* 41 (1958) 40.
- [34] H.R. Kirchmayr, C.A. Poldy, *J. Magn. Magn. Mater.* 8 (1978) 1.
- [35] Z. Li, G. Chen, X. Tian, Y. Li, *Mater. Res. Bull.* 43 (2008) 1781.
- [36] A.V. Bune, V.M. Fridkin, S. Ducharme, L.M. Blinov, S.P. Palto, A.V. Sorokin, S.G. Yudin, A. Zlatkin, *Nature* 391 (1998) 874.
- [37] W.L. Zhong, B. Jiang, P.L. Zhang, J.M. Ma, H.M. Cheng, Z.H. Yang, L.X. Li, *J. Phys. Condens. Matter* 5 (1993) 2619.
- [38] E.D. Mishina, K.A. Vorotilov, V.A. Vasiliev, A.S. Sigov, N. Ota, S. Nakabayashi, *J. Exp. Theor. Phys.* 122 (2002) 582.
- [39] G. Geneste, E. Bousquet, J. Junquera, P. Ghosez, *Appl. Phys. Lett.* 88 (2006) 112906.
- [40] J. Hao, Z. Xu, R. Chu, W. Li, G. Li, Q. Yin, *J. Alloys Compd.* 484 (2009) 233.
- [41] J. Wu, J. Wang, *J. Appl. Phys.* 105 (2009) 124107.
- [42] R. Font, O. Raymond, E. Martinez, J. Portelles, J.M. Siqueiros, *J. Appl. Phys.* 105 (2009) 114110.